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Editors

# Studies in Phase Space Analysis with Applications to PDEs



# **Progress in Nonlinear Differential Equations and Their Applications**

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# Studies in Phase Space Analysis with Applications to PDEs

 Birkhäuser

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*Mais où sont les neiges d'antan?*

F. Villon



# Preface

This volume is a collection of papers mainly concerning *phase space analysis*, or *microlocal analysis*, and its applications to the theory of *partial differential equations* (PDEs).

A number of papers composing this volume, all written by leading experts in their respective fields, are expanded version of talks given at a meeting held in September 2011 at the *Centro Residenziale Universitario (CEUB)* of Bertinoro, on the hills surrounding Cesena, Italy.

The Bertinoro workshop was the occasion to fix the state of the art in many different aspects of phase space analysis. In fact the results collected here concern general theory of pseudodifferential operators, Hardy-type inequalities, linear and non-linear hyperbolic equations and systems, water-waves equations, Euler–Poisson and Navier–Stokes equations, Schrödinger equations and heat and parabolic equations.

We would like to seize this occasion to thank all the contributors as well as the people who took part in the workshop.

A number of institutions have made it possible to hold the Bertinoro workshop through their financial support. We’d like to list them here: the Italian Ministero dell’Istruzione, dell’Università e della Ricerca by means of the PRIN 2008 project “Phase Space Analysis of PDE’s”, the Istituto Nazionale di Alta Matematica, the University of Bordeaux 1, the University of Pisa and the Fondazione Cassa di Risparmio di Cesena. We thank all of them for their generosity.

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# Chapter 1

## The Water-Wave Equations: From Zakharov to Euler

Thomas Alazard, Nicolas Burq, and Claude Zuily

**Abstract** Starting from the Zakharov/Craig–Sulem formulation of the water-wave equations, we prove that one can define a pressure term and hence obtain a solution of the classical Euler equations. It is proved that these results hold in rough domains, under minimal assumptions on the regularity to ensure, in terms of Sobolev spaces, that the solutions are  $C^1$ .

**Key words:** Cauchy theory, Euler equations, Water-wave system

**Mathematics Subject Classification:** 35B65, 35B30, 35B60, 35J75, 35S05.

### 1.1 Introduction

We study the dynamics of an incompressible layer of inviscid liquid, having constant density, occupying a fluid domain with a free surface.

We begin by describing the fluid domain. Hereafter,  $d \geq 1$ ,  $t$  denotes the time variable, and  $x \in \mathbf{R}^d$  and  $y \in \mathbf{R}$  denote the horizontal and vertical variables. We work in a fluid domain with free boundary of the form

$$\Omega = \{ (t, x, y) \in (0, T) \times \mathbf{R}^d \times \mathbf{R} : (x, y) \in \Omega(t) \},$$

where  $\Omega(t)$  is the  $d + 1$ -dimensional domain located between two hypersurfaces: a free surface denoted by  $\Sigma(t)$  which will be supposed to be a graph and a fixed

---

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bottom  $\Gamma$ . For each time  $t$ , one has

$$\Omega(t) = \{(x, y) \in \mathcal{O} : y < \eta(t, x)\},$$

where  $\mathcal{O}$  is a given open connected domain and where  $\eta$  is the free surface elevation. We denote by  $\Sigma$  the free surface

$$\Sigma = \{(t, x, y) : t \in (0, T), (x, y) \in \Sigma(t)\},$$

where  $\Sigma(t) = \{(x, y) \in \mathbf{R}^d \times \mathbf{R} : y = \eta(t, x)\}$ , and we set  $\Gamma = \partial\Omega(t) \setminus \Sigma(t)$ .

Notice that  $\Gamma$  does not depend on time. Two classical examples are the case of infinite depth ( $\mathcal{O} = \mathbf{R}^{d+1}$  so that  $\Gamma = \emptyset$ ) and the case where the bottom is the graph of a function (this corresponds to the case  $\mathcal{O} = \{(x, y) \in \mathbf{R}^d \times \mathbf{R} : y > b(x)\}$  for some given function  $b$ ).

We introduce now a condition which ensures that, at time  $t$ , there exists a fixed strip separating the free surface from the bottom:

$$(H_t) : \quad \exists h > 0 : \quad \Gamma \subset \{(x, y) \in \mathbf{R}^d \times \mathbf{R} : y < \eta(t, x) - h\}. \quad (1.1)$$

No regularity assumption will be made on the bottom  $\Gamma$ .

## The Incompressible Euler Equation with Free Surface

Hereafter, we use the following notations:

$$\nabla = (\partial_{x_i})_{1 \leq i \leq d}, \quad \nabla_{x,y} = (\nabla, \partial_y), \quad \Delta = \sum_{1 \leq i \leq d} \partial_{x_i}^2, \quad \Delta_{x,y} = \Delta + \partial_y^2.$$

The Eulerian velocity field  $v: \Omega \rightarrow \mathbf{R}^{d+1}$  solves the incompressible Euler equation

$$\partial_t v + v \cdot \nabla_{x,y} v + \nabla_{x,y} P = -g e_y, \quad \operatorname{div}_{x,y} v = 0 \quad \text{in } \Omega,$$

where  $g$  is the acceleration due to gravity ( $g > 0$ ) and  $P$  is the pressure. The problem is then given by three boundary conditions:

- A kinematic condition (which states that the free surface moves with the fluid)

$$\partial_t \eta = \sqrt{1 + |\nabla \eta|^2} (v \cdot n) \quad \text{on } \Sigma, \quad (1.2)$$

where  $n$  is the unit exterior normal to  $\Omega(t)$

- A dynamic condition (that expresses a balance of forces across the free surface)

$$P = 0 \quad \text{on } \Sigma \quad (1.3)$$

- The ‘‘solid wall’’ boundary condition at the bottom  $\Gamma$

$$v \cdot \nu = 0, \quad (1.4)$$

where  $\nu$  is the normal vector to  $\Gamma$  whenever it exists. In the case of arbitrary bottom, this condition will be implicit and contained in a variational formulation.

### The Zakharov/Craig–Sulem Formulation

A popular form of the water-wave system is given by the Zakharov/Craig–Sulem formulation. This is an elegant formulation of the water-wave equations where all the unknowns are evaluated at the free surface only. Let us recall the derivation of this system.

Assume, furthermore, that the motion of the liquid is irrotational. The velocity field  $\nu$  is therefore given by  $\nu = \nabla_{x,y}\Phi$  for some velocity potential  $\Phi: \Omega \rightarrow \mathbf{R}$  satisfying

$$\Delta_{x,y}\Phi = 0 \quad \text{in } \Omega, \quad \partial_\nu\Phi = 0 \quad \text{on } \Gamma,$$

and the Bernoulli equation

$$\partial_t\Phi + \frac{1}{2}|\nabla_{x,y}\Phi|^2 + P + gy = 0 \quad \text{in } \Omega. \quad (1.5)$$

Following Zakharov [7], introduce the trace of the potential on the free surface:

$$\psi(t, x) = \Phi(t, x, \eta(t, x)).$$

Notice that since  $\Phi$  is harmonic,  $\eta$  and  $\Psi$  fully determines  $\Phi$ . Craig and Sulem (see [3]) observe that one can form a system of two evolution equations for  $\eta$  and  $\psi$ . To do so, they introduce the Dirichlet-Neumann operator  $G(\eta)$  that relates  $\psi$  to the normal derivative  $\partial_n\Phi$  of the potential by

$$\begin{aligned} (G(\eta)\psi)(t, x) &= \sqrt{1 + |\nabla\eta|^2} \partial_n\Phi|_{y=\eta(t, x)} \\ &= (\partial_y\Phi)(t, x, \eta(t, x)) - \nabla_x\eta(t, x) \cdot (\nabla_x\Phi)(t, x, \eta(t, x)). \end{aligned}$$

(For the case with a rough bottom, we recall the precise construction later on). Directly from this definition, one has

$$\partial_t\eta = G(\eta)\psi. \quad (1.6)$$

It is proved in [3] (see also the computations in §1.3.6) that the condition  $P = 0$  on the free surface implies that

$$\partial_t\psi + g\eta + \frac{1}{2}|\nabla\psi|^2 - \frac{1}{2} \frac{(\nabla\eta \cdot \nabla\psi + G(\eta)\psi)^2}{1 + |\nabla\eta|^2} = 0. \quad (1.7)$$

The system (1.6) and (1.7) is in Hamiltonian form (see [3, 7]), where the Hamiltonian is given by

$$\mathcal{H} = \frac{1}{2} \int_{\mathbf{R}^d} \psi G(\eta)\psi + g\eta^2 dx.$$

The problem to be considered here is that of the equivalence of the previous two formulations of the water-wave problem. Assume that the Zakharov/Craig–Sulem system has been solved. Namely, assume that, for some  $r > 1 + d/2$ ,  $(\eta, \psi) \in C^0(I, H^r(\mathbf{R}^d) \times H^r(\mathbf{R}^d))$  solves (1.6)–(1.7). We would like to show that we have indeed solved the initial system of Euler’s equation with free boundary. In particular we have to define the pressure which does not appear in the above system (1.6)–(1.7). To do so, we set

$$B = \frac{\nabla\eta \cdot \nabla\psi + G(\eta)\psi}{1 + |\nabla\eta|^2}, \quad V = \nabla\psi - B\nabla\eta.$$

Then  $B$  and  $V$  belong to the space  $C^0(I, H^{\frac{1}{2}}(\mathbf{R}^d))$ . It follows from [1] that (for fixed  $t$ ) one can define unique variational solutions to the problems

$$\Delta_{x,y}\Phi = 0 \quad \text{in } \Omega, \quad \Phi|_{\Sigma} = \psi, \quad \partial_\nu\Phi = 0 \quad \text{on } \Gamma.$$

$$\Delta_{x,y}Q = 0 \quad \text{in } \Omega, \quad Q|_{\Sigma} = g\eta + \frac{1}{2}(B^2 + |V|^2), \quad \partial_\nu Q = 0 \quad \text{on } \Gamma.$$

Then we shall define  $P \in \mathcal{D}'(\Omega)$  by

$$P := Q - gy - \frac{1}{2}|\nabla_{x,y}\Phi|^2$$

and we shall show firstly that  $P$  has a trace on  $\Sigma$  which is equal to 0 and secondly that  $Q = -\partial_t\Phi$  which will show, according to (1.5), that we have indeed solved Bernoulli’s (and therefore Euler’s) equation.

These assertions are not straightforward because we are working with solutions of low regularity and we consider general bottoms (namely no regularity assumption is assumed on the bottom). Indeed, the analysis would have been much easier for  $r > 2 + d/2$  and a flat bottom.

## 1.2 Low Regularity Cauchy Theory

Since we are interested in low regularity solutions, we begin by recalling the well-posedness results proved in [2]. These results clarify the Cauchy theory of the water-wave equations as well in terms of regularity indexes for the initial conditions as for the smoothness of the bottom of the domain (namely no regularity assumption is assumed on the bottom).

Recall that the Zakharov/Craig–Sulem system reads

$$\begin{cases} \partial_t\eta - G(\eta)\psi = 0, \\ \partial_t\psi + g\eta + \frac{1}{2}|\nabla\psi|^2 - \frac{1}{2} \frac{(\nabla\eta \cdot \nabla\psi + G(\eta)\psi)^2}{1 + |\nabla\eta|^2} = 0. \end{cases} \quad (1.8)$$

It is useful to introduce the vertical and horizontal components of the velocity

$$B := (v_y)|_{y=\eta} = (\partial_y \Phi)|_{y=\eta}, \quad V := (v_x)|_{y=\eta} = (\nabla_x \Phi)|_{y=\eta}.$$

These can be defined in terms of  $\eta$  and  $\psi$  by means of the formulas

$$B = \frac{\nabla \eta \cdot \nabla \psi + G(\eta) \psi}{1 + |\nabla \eta|^2}, \quad V = \nabla \psi - B \nabla \eta. \quad (1.9)$$

Also, recall that the Taylor coefficient  $a = -\partial_y P|_{\Sigma}$  can be defined in terms of  $\eta, V, B,$  and  $\psi$  only (see §4.3.1 in [5]).

In [2] we proved the following results about low regularity solutions. We refer to the introduction of [2, 4] for references and a short historical survey of the background of this problem.

**Theorem 1.1 ([2]).** *Let  $d \geq 1$  and  $s > 1 + d/2$  and consider an initial data  $(\eta_0, \psi_0)$  such that:*

- (i)  $\eta_0 \in H^{s+\frac{1}{2}}(\mathbf{R}^d), \quad \psi_0 \in H^{s+\frac{1}{2}}(\mathbf{R}^d), \quad V_0 \in H^s(\mathbf{R}^d), \quad B_0 \in H^s(\mathbf{R}^d).$
- (ii) *The condition  $(H_0)$  in (1.1) holds initially for  $t = 0$ .*
- (iii) *There exists a positive constant  $c$  such that, for all  $x$  in  $\mathbf{R}^d, a_0(x) \geq c$ .*

*Then there exists  $T > 0$  such that the Cauchy problem for (1.8) with initial data  $(\eta_0, \psi_0)$  has a unique solution:*

$$(\eta, \psi) \in C^0([0, T], H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d)),$$

*such that*

1.  $(V, B) \in C^0([0, T], H^s(\mathbf{R}^d) \times H^s(\mathbf{R}^d)).$
2. *The condition  $(H_t)$  in (1.1) holds for  $t \in [0, T]$  with  $h$  replaced by  $h/2$ .*
3.  $a(t, x) \geq c/2,$  for all  $(t, x)$  in  $[0, T] \times \mathbf{R}^d$ .

In a forthcoming paper we shall prove the following result.

**Theorem 1.2.** *Assume  $\Gamma = \emptyset$ . Let  $d = 2$  and  $s > 1 + \frac{d}{2} - \frac{1}{12}$  and consider an initial data  $(\eta_0, \psi_0)$  such that*

$$\eta_0 \in H^{s+\frac{1}{2}}(\mathbf{R}^d), \quad \psi_0 \in H^{s+\frac{1}{2}}(\mathbf{R}^d), \quad V_0 \in H^s(\mathbf{R}^d), \quad B_0 \in H^s(\mathbf{R}^d).$$

*Then there exists  $T > 0$  such that the Cauchy problem for (1.8) with initial data  $(\eta_0, \psi_0)$  has a solution  $(\eta, \psi)$  such that*

$$(\eta, \psi, V, B) \in C^0([0, T]; H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^s(\mathbf{R}^d) \times H^s(\mathbf{R}^d)).$$

*Remark 1.1.* (i) For the sake of simplicity we stated Theorem 1.2 in dimension  $d = 2$  (recall that  $d$  is the dimension of the interface). One can prove such a result in any dimension  $d \geq 2$ , the number  $1/12$  being replaced by an index depending on  $d$ .

(ii) Notice that in infinite depth ( $\Gamma = \emptyset$ ) the Taylor condition (which is assumption (iii) in Theorem 1.1) is always satisfied as proved by Wu ([6]).

Now having solved the system (1.8) in  $(\eta, \psi)$ , we have to show that we have indeed solved the initial system in  $(\eta, v)$ . This is the purpose of the following section.

There is one point that should be emphasized concerning the regularity. Below we consider solutions  $(\eta, \psi)$  of (1.8) such that

$$(\eta, \psi) \in C^0([0, T]; H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d)),$$

with the only assumption that  $s > \frac{1}{2} + \frac{d}{2}$  (and the assumption that there exists  $h > 0$  such that the condition  $(H_t)$  in (1.1) holds for  $t \in [0, T]$ ). Consequently, the result proved in this note applies to the settings considered in the above theorems.

## 1.3 From Zakharov to Euler

### 1.3.1 The Variational Theory

In this paragraph the time is fixed so we will skip it and work in a fixed domain  $\Omega$  whose top boundary  $\Sigma$  is Lipschitz, i.e.  $\eta \in W^{1,\infty}(\mathbf{R}^d)$ .

We recall here the variational theory, developed in [1], allowing us to solve the following problem in the case of arbitrary bottom,

$$\Delta \Phi = 0 \quad \text{in } \Omega, \quad \Phi|_{\Sigma} = \psi, \quad \frac{\partial \Phi}{\partial \nu}|_{\Gamma} = 0. \quad (1.10)$$

Notice that  $\Omega$  is not necessarily bounded below. We proceed as follows.

Denote by  $\mathcal{D}$  the space of functions  $u \in C^\infty(\Omega)$  such that  $\nabla_{x,y} u \in L^2(\Omega)$ , and let  $\mathcal{D}_0$  be the subspace of functions  $u \in \mathcal{D}$  such that  $u$  vanishes near the top boundary  $\Sigma$ .

**Lemma 1.1** (see Prop 2.2 in [1]). *There exist a positive weight  $g \in L^\infty_{loc}(\Omega)$  equal to 1 near the top boundary  $\Sigma$  of  $\Omega$  and  $C > 0$  such that for all  $u \in \mathcal{D}_0$ ,*

$$\iint_{\Omega} g(x,y) |u(x,y)|^2 dx dy \leq C \iint_{\Omega} |\nabla_{x,y} u(x,y)|^2 dx dy. \quad (1.11)$$

Using this lemma one can prove the following result.

**Proposition 1.1** (see page 422 in [1]). *Denote by  $H^{1,0}(\Omega)$  the space of functions  $u$  on  $\Omega$  such that there exists a sequence  $(u_n) \subset \mathcal{D}_0$  such that*

$$\nabla_{x,y} u_n \rightarrow \nabla_{x,y} u \quad \text{in } L^2(\Omega), \quad u_n \rightarrow u \quad \text{in } L^2(\Omega, g dx dy),$$

*endowed with the scalar product*

$$(u, v)_{H^{1,0}(\Omega)} = (\nabla_x u, \nabla_x v)_{L^2(\Omega)} + (\partial_y u, \partial_y v)_{L^2(\Omega)}.$$

*Then  $H^{1,0}(\Omega)$  is a Hilbert space and (1.11) holds for  $u \in H^{1,0}(\Omega)$ .*



Let  $\underline{\psi} \in H^{\frac{1}{2}}(\mathbf{R}^d)$ . One can construct (see below after (1.21))  $\underline{\psi} \in H^1(\Omega)$  such that

$$\text{supp } \underline{\psi} \subset \{(x, y) : \eta(t, x) - h \leq y \leq \eta(x)\}, \quad \underline{\psi}|_{\Sigma} = \psi.$$

Using Proposition 1.1 we deduce that there exists a unique  $u \in H^{1,0}(\Omega)$  such that, for all  $\theta \in H^{1,0}(\Omega)$ ,

$$\iint_{\Omega} \nabla_{x,y} u(x, y) \cdot \nabla_{x,y} \theta(x, y) dx dy = - \iint_{\Omega} \nabla_{x,y} \underline{\psi}(x, y) \cdot \nabla_{x,y} \theta(x, y) dx dy.$$

Then to solve the problem (1.10) we set  $\Phi = u + \underline{\psi}$ .

*Remark 1.2.* As for the usual Neumann problem the meaning of the third condition in (1.10) is included in the definition of the space  $H^{1,0}(\Omega)$ . It can be written as in (1.10) if the bottom  $\Gamma$  is sufficiently smooth.

### 1.3.2 The Main Result

Let us assume that the Zakharov system (1.8) has been solved on  $I = (0, T)$ , which means that we have found, for  $s > \frac{1}{2} + \frac{d}{2}$ , a solution

$$(\eta, \psi) \in C^0(\bar{I}, H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d)),$$

of the system

$$\begin{cases} \partial_t \eta = G(\eta) \psi, \\ \partial \psi = -g\eta - \frac{1}{2} |\nabla \psi|^2 + \frac{1}{2} \frac{(\nabla \psi \cdot \nabla \eta + G(\eta) \psi)^2}{1 + |\nabla \eta|^2}. \end{cases} \quad (1.12)$$

Let  $B, V$  be defined by (1.9). Then  $(B, V) \in C^0(I, H^{s-\frac{1}{2}}(\mathbf{R}^d) \times H^{s-\frac{1}{2}}(\mathbf{R}^d))$ .

The above variational theory shows that one can solve (for fixed  $t$ ) the problem

$$\Delta_{x,y} Q = 0 \quad \text{in } \Omega, \quad Q|_{\Sigma} = g\eta + \frac{1}{2}(B^2 + |V|^2) \in H^{\frac{1}{2}}(\mathbf{R}^d). \quad (1.13)$$

Here is the main result of this article.

**Theorem 1.3.** *Let  $\Phi$  and  $Q$  be the variational solutions of the problems (1.10) and (1.13). Set  $P = Q - g\eta - \frac{1}{2} |\nabla_{x,y} \Phi|^2$ . Then  $v := \nabla_{x,y} \Phi$  satisfies the Euler system*

$$\partial_t v + (v \cdot \nabla_{x,y}) v + \nabla_{x,y} P = -g e_y \quad \text{in } \Omega,$$

together with the conditions

$$\begin{cases} \operatorname{div}_{x,y} v = 0, & \operatorname{curl}_{x,y} v = 0 & \text{in } \Omega, \\ \partial_t \eta = (1 + |\nabla \eta|^2)^{\frac{1}{2}} (v \cdot n) & & \text{on } \Sigma, \\ P = 0 & & \text{on } \Sigma. \end{cases} \quad (1.14)$$

The rest of the paper is devoted to the proof of this result. We proceed in several steps.

### 1.3.3 Straightening the Free Boundary

First of all if condition  $(H_t)$  is satisfied on  $I$ , for  $T$  small enough, one can find  $\eta_* \in L^\infty(\mathbf{R}^d)$  independent of  $t$  such that

$$\begin{cases} (i) & \nabla_x \eta_* \in H^\infty(\mathbf{R}^d), \quad \|\nabla_x \eta_*\|_{L^\infty(\mathbf{R}^d)} \leq C \|\eta\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d))}, \\ (ii) & \eta(t, x) - h \leq \eta_*(x) \leq \eta(t, x) - \frac{h}{2}, \quad \forall (t, x) \in I \times \mathbf{R}^d, \\ (iii) & \Gamma \subset \{(x, y) \in \mathcal{O} : y < \eta_*(x)\}. \end{cases} \quad (1.15)$$

Indeed using the first equation in (1.12) we have

$$\begin{aligned} \|\eta(t, \cdot) - \eta_0\|_{L^\infty(\mathbf{R}^d)} &\leq \int_0^t \|G(\eta) \psi(\sigma, \cdot)\|_{H^{s-\frac{1}{2}}(\mathbf{R}^d)} d\sigma \\ &\leq TC \left( \|\eta, \psi\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d))} \right). \end{aligned}$$

Therefore taking  $T$  small enough we make  $\|\eta(t, \cdot) - \eta_0\|_{L^\infty(\mathbf{R}^d)}$  as small as we want.

Then we take  $\eta_*(x) = -\frac{2h}{3} + e^{-\delta|D_x|} \eta_0$  and writing

$$\eta_*(x) = -\frac{2h}{3} + \eta(t, x) - (\eta(t, x) - \eta_0(x)) + (e^{-\delta|D_x|} \eta_0 - \eta_0(x)),$$

we obtain (1.15).

In what follows we shall set

$$\begin{cases} \Omega_1(t) = \{(x, y) : x \in \mathbf{R}^d, \eta_*(x) < y < \eta(t, x)\}, \\ \Omega_1 = \{(t, x, y) : t \in I, (x, y) \in \Omega_1(t)\}, & \Omega_2 = \{(x, y) \in \mathcal{O} : y \leq \eta_*(x)\}, \\ \tilde{\Omega}_1 = \{(x, z) : x \in \mathbf{R}^d, z \in (-1, 0)\}, \\ \tilde{\Omega}_2 = \{(x, z) \in \mathbf{R}^d \times (-\infty, -1] : (x, z+1 + \eta_*(x)) \in \Omega_2\} \\ \tilde{\Omega} = \tilde{\Omega}_1 \cup \tilde{\Omega}_2 \end{cases} \quad (1.16)$$

Following Lannes ([4]), for  $t \in I$  consider the map  $(x, z) \mapsto (x, \rho(t, x, z))$  from  $\tilde{\Omega}$  to  $\mathbf{R}^{d+1}$  defined by

$$\begin{cases} \rho(t, x, z) = (1+z)e^{\delta z \langle D_x \rangle} \eta(t, x) - z \eta_*(x) & \text{if } (x, z) \in \tilde{\Omega}_1 \\ \rho(t, x, z) = z+1 + \eta_*(x) & \text{if } (x, z) \in \tilde{\Omega}_2. \end{cases} \quad (1.17)$$

where  $\delta$  is chosen such that

$$\delta \|\eta\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d))} := \delta_0 \ll 1.$$

Notice that since  $s > \frac{1}{2} + \frac{d}{2}$ , taking  $\delta$  small enough and using (1.15) (i) and, (ii), we obtain the estimates

$$\begin{aligned} (i) \quad \partial_z \rho(t, x, z) &\geq \min\left(\frac{h}{3}, 1\right) \quad \forall (t, x, z) \in I \times \tilde{\Omega}, \\ (ii) \quad \|\nabla_{x,z} \rho\|_{L^\infty(I \times \tilde{\Omega})} &\leq C(1 + \|\eta\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d))}). \end{aligned} \quad (1.18)$$

It follows from (1.18) (i) that the map  $(t, x, z) \mapsto (t, x, \rho(t, x, z))$  is a diffeomorphism from  $I \times \tilde{\Omega}$  to  $\Omega$  which is of class  $W^{1,\infty}$ .

We denote by  $\kappa$  the inverse map of  $\rho$ :

$$\begin{aligned} (t, x, z) \in I \times \tilde{\Omega}, (t, x, \rho(t, x, z)) &= (t, x, y) \\ \iff (t, x, z) &= (t, x, \kappa(t, x, y)), (t, x, y) \in \Omega. \end{aligned} \quad (1.19)$$

### 1.3.4 The Dirichlet-Neumann Operator

Let  $\Phi$  be the variational solution described above (with fixed  $t$ ) of the problem

$$\begin{cases} \Delta_{x,y} \Phi = 0 & \text{in } \Omega(t), \\ \Phi|_{\Sigma(t)} = \psi(t, \cdot), \\ \partial_\nu \Phi|_\Gamma = 0. \end{cases} \quad (1.20)$$

Let us recall that

$$\Phi = u + \underline{\psi} \quad (1.21)$$

where  $u \in H^{1,0}(\Omega(t))$  and  $\underline{\psi}$  is an extension of  $\psi$  to  $\Omega(t)$ .

Here is a construction of  $\underline{\psi}$ . Let  $\chi \in C^\infty(\mathbf{R})$ ,  $\chi(a) = 0$  if  $a \leq -1$ ,  $\chi(a) = 1$  if  $a \geq -\frac{1}{2}$ . Let  $\tilde{\psi}(t, x, z) = \chi(z) e^{z(D_x)} \psi(t, x)$  for  $z \leq 0$ . It is classical that  $\tilde{\psi} \in L^\infty(I, H^1(\tilde{\Omega}))$  if  $\psi \in L^\infty(I, H^{\frac{1}{2}}(\mathbf{R}^d))$  and

$$\|\tilde{\psi}\|_{L^\infty(I, H^1(\tilde{\Omega}))} \leq C \|\psi\|_{L^\infty(I, H^{\frac{1}{2}}(\mathbf{R}^d))}.$$

Then we set

$$\underline{\psi}(t, x, y) = \tilde{\psi}(t, x, \kappa(t, x, y)). \quad (1.22)$$

Since  $\eta \in C^0(I, W^{1,\infty}(\mathbf{R}^d))$  we have  $\underline{\psi}(t, \cdot) \in H^1(\Omega(t))$ ,  $\underline{\psi}|_{\Sigma(t)} = \psi$ , and

$$\|\underline{\psi}(t, \cdot)\|_{H^1(\Omega(t))} \leq C(\|\eta\|_{L^\infty(I, W^{1,\infty}(\mathbf{R}^d))}) \|\psi\|_{L^\infty(I, H^{\frac{1}{2}}(\mathbf{R}^d))}.$$

Then we define the Dirichlet-Neumann operator by

$$\begin{aligned} G(\eta)\psi(t,x) &= \sqrt{1 + |\nabla\eta|^2} \partial_n \Phi|_{\Sigma} \\ &= (\partial_y \Phi)(x, \eta(t,x)) - \nabla_x \eta(t,x) \cdot (\nabla_x \Phi)(t,x, \eta(t,x)). \end{aligned} \quad (1.23)$$

It has been shown in [2] (see Sect. 3) that  $G(\eta)\psi$  is well defined in  $C^0(\bar{I}, H^{-\frac{1}{2}}(\mathbf{R}^d))$  if  $\eta \in C^0(\bar{I}, W^{1,\infty}(\mathbf{R}^d))$  and  $\psi \in C^0(\bar{I}, H^{\frac{1}{2}}(\mathbf{R}^d))$ .

*Remark 1.3.* Recall that we have set

$$\Omega(t) = \{(x,y) \in \mathcal{O} : y < \eta(t,x)\}, \quad \Omega = \{(t,x,y) : t \in I, (x,y) \in \Omega(t)\}. \quad (1.24)$$

For a function  $f \in L^1_{loc}(\Omega)$  if  $\partial_t f$  denotes its derivative in the sense of distributions, we have

$$\langle \partial_t f, \varphi \rangle = \lim_{\varepsilon \rightarrow 0} \left\langle \frac{f(\cdot + \varepsilon, \cdot, \cdot) - f(\cdot, \cdot, \cdot)}{\varepsilon}, \varphi \right\rangle, \quad \forall \varphi \in C_0^\infty(\Omega). \quad (1.25)$$

This point should be clarified due to the particular form of the set  $\Omega$  since we have to show that if  $(t,x,y) \in \text{supp } \varphi = K$ , then  $(t + \varepsilon, x, y) \in \Omega$  for  $\varepsilon$  sufficiently small independently of the point  $(t,x,y)$ . This is true. Indeed if  $(t,x,y) \in K$ , there exists a fixed  $\delta > 0$  (depending only on  $K, \eta$ ) such that  $y \leq \eta(t,x) - \delta$ . Since by (1.12)

$$|\eta(t + \varepsilon, x) - \eta(t, x)| \leq \varepsilon \|G(\eta)\psi\|_{L^\infty(I \times \mathbf{R}^d)} \leq \varepsilon C$$

where  $C = C(\|(\eta, \psi)\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d))})$ , we have if  $\varepsilon < \frac{\delta}{C}$ ,

$$y - \eta(t + \varepsilon, x) = y - \eta(t, x) + \eta(t, x) - \eta(t + \varepsilon, x) \leq -\delta + \varepsilon C < 0.$$

Notice that since  $\eta \in C^0(\bar{I}, H^{s+\frac{1}{2}}(\mathbf{R}^d))$ ,  $\partial_t \eta = G(\eta)\psi \in C^0(\bar{I}, H^{s-\frac{1}{2}}(\mathbf{R}^d))$ , and  $s > \frac{1}{2} + \frac{d}{2}$ , we have  $\rho \in W^{1,\infty}(I \times \tilde{\Omega})$ .

The main step in the proof of Theorem 1.3 is the following.

**Proposition 1.2.** *Let  $\Phi$  be defined by (1.20) and  $Q \in H^{1,0}(\Omega(t))$  by (1.13). Then for all  $t \in I$*

- (i)  $\partial_t \Phi(t, \cdot) \in H^{1,0}(\Omega(t))$ .
- (ii)  $\partial_t \Phi = -Q$  in  $\mathcal{D}'(\Omega)$ .

This result will be proved in §1.3.6

### 1.3.5 Preliminaries

If  $f$  is a function defined on  $\Omega$ , we shall denote by  $\tilde{f}$  its image by the diffeomorphism  $(t,x,z) \mapsto (t,x,\rho(t,x,z))$ . Thus we have

$$\tilde{f}(t, x, z) = f(t, x, \rho(t, x, z)) \Leftrightarrow f(t, x, y) = \tilde{f}(t, x, \kappa(t, x, y)). \quad (1.26)$$

Formally we have the following equalities for  $(t, x, y) = (t, x, \rho(t, x, z)) \in \Omega$  and  $\nabla = \nabla_x$

$$\begin{cases} \partial_y f(t, x, y) = \frac{1}{\partial_z \rho} \partial_z \tilde{f}(t, x, z) \Leftrightarrow \partial_z \tilde{f}(t, x, z) = \partial_z \rho(t, x, \kappa(t, x, y)) \partial_y f(t, x, y), \\ \nabla f(t, x, y) = (\nabla \tilde{f} - \frac{\nabla \rho}{\partial_z \rho} \partial_z \tilde{f})(t, x, z) \Leftrightarrow \nabla \tilde{f}(t, x, z) = (\nabla f + \nabla \rho \partial_y f)(t, x, y), \\ \partial_t f(t, x, y) = (\partial_t \tilde{f} + \partial_t \kappa(t, x, y) \partial_z \tilde{f})(t, x, \kappa(t, x, y)). \end{cases} \quad (1.27)$$

We shall set in what follows

$$\Lambda_1 = \frac{1}{\partial_z \rho} \partial_z, \quad \Lambda_2 = \nabla_x - \frac{\nabla_x \rho}{\partial_z \rho} \partial_z \quad (1.28)$$

Eventually recall that if  $u$  is the function defined by (1.21), we have

$$\iint_{\Omega(t)} \nabla_{x,y} u(t, x, y) \cdot \nabla_{x,y} \theta(x, y) dx dy = - \iint_{\Omega(t)} \nabla_{x,y} \underline{\psi}(t, x, y) \cdot \nabla_{x,y} \theta(x, y) dx dy \quad (1.29)$$

for all  $\theta \in H^{1,0}(\Omega(t))$  which implies that for  $t \in I$ ,

$$\|\nabla_{x,y} u(t, \cdot)\|_{L^2(\Omega(t))} \leq C(\|\eta\|_{L^\infty(I, W^{1,\infty}(\mathbf{R}^d))}) \|\psi\|_{L^\infty(I, H^{\frac{1}{2}}(\mathbf{R}^d))}. \quad (1.30)$$

Let  $u$  be defined by (1.21). Since  $(\eta, \psi) \in C^0(\bar{I}, H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d))$ , the elliptic regularity theorem proved in [2] (see Theorem 3.16) shows that

$$\partial_z \tilde{u}, \nabla_x \tilde{u} \in C_z^0([-1, 0], H^{s-\frac{1}{2}}(\mathbf{R}^d)) \subset C^0([-1, 0] \times \mathbf{R}^d),$$

since  $s - \frac{1}{2} > \frac{d}{2}$ .

It follows from (1.27) that  $\partial_y u$  and  $\nabla_x u$  have a trace on  $\Sigma$  and

$$\partial_y u|_\Sigma = \frac{1}{\partial_z \rho(t, x, 0)} \partial_z \tilde{u}(t, x, 0), \quad \nabla_x u|_\Sigma = (\nabla_x \tilde{u} - \frac{\nabla_x \eta}{\partial_z \rho(t, x, 0)} \partial_z \tilde{u})(t, x, 0).$$

Since  $\tilde{u}(t, x, 0) = 0$ , it follows that

$$\nabla_x u|_\Sigma + (\nabla_x \eta) \partial_y u|_\Sigma = 0$$

from which we deduce, since  $\Phi = u + \underline{\psi}$ ,

$$\nabla_x \Phi|_\Sigma + (\nabla_x \eta) \partial_y \Phi|_\Sigma = \nabla_x \psi. \quad (1.31)$$

On the other hand one has

$$G(\eta) \psi = (\partial_y \Phi - \nabla_x \eta \cdot \nabla_x \Phi)|_\Sigma. \quad (1.32)$$

It follows from (1.31), (1.32), and (1.9) that we have

$$\nabla_x \Phi|_{\Sigma} = V, \quad \partial_y \Phi|_{\Sigma} = B. \quad (1.33)$$

According to (1.13),  $P = Q - g\eta - \frac{1}{2}|\nabla_{x,y}\Phi|^2$  has a trace on  $\Sigma$  and  $P|_{\Sigma} = 0$ .

### 1.3.6 The Regularity Results

The main steps in the proof of Proposition 1.2 are the following.

**Lemma 1.2.** *Let  $\tilde{u}$  be defined by (1.26) and  $\kappa$  by (1.19). Then for all  $t_0 \in I$  the function  $(x, y) \mapsto U(t_0, x, y) := \partial_t \tilde{u}(t_0, x, \kappa(t_0, x, y))$  belongs to  $H^{1,0}(\Omega(t_0))$ . Moreover there exists a function  $\mathcal{F} : \mathbf{R}^+ \rightarrow \mathbf{R}^+$  such that*

$$\sup_{t \in I} \iint_{\Omega(t)} |\nabla_{x,y} U(t, x, y)|^2 dx dy \leq \mathcal{F}(\|(\eta, \psi)\|_{L^\infty(I, H^{s+\frac{1}{2}}(\mathbf{R}^d) \times H^{s+\frac{1}{2}}(\mathbf{R}^d))}).$$

**Lemma 1.3.** *In the sense of distributions on  $\Omega$  we have the chain rule*

$$\partial_t u(t, x, y) = \partial_t \tilde{u}(t, x, \kappa(t, x, y)) + \partial_t \kappa(t, x, y) \partial_z \tilde{u}(t, x, \kappa(t, x, y)).$$

These lemmas are proved in the next paragraph.

*Proof (of Proposition 1.2).* According to (1.21) and Lemma 1.3 we have

$$\partial_t \Phi(t, x, y) = \partial_t \tilde{u}(t, x, \kappa(t, x, y)) + \underline{w}(t, x, y) \quad (1.34)$$

where

$$\underline{w}(t, x, y) = \partial_t \kappa(t, x, y) \partial_z \tilde{u}(t, x, \kappa(t, x, y)) + \partial_t \underline{\psi}(t, x, y).$$

According to Lemma 1.2 the first term in the right-hand side of (1.34) belongs to  $H^{1,0}(\Omega(t))$ . Denoting by  $\tilde{w}$  the image of  $\underline{w}$ , if we show that

$$\begin{cases} (i) & \tilde{w} \in H^1(\mathbf{R}^d \times \mathbf{R}), \\ (ii) & \text{supp } \tilde{w} \subset \{(x, z) \in \mathbf{R}^d \times (-1, 0)\} \\ (iii) & w|_{\Sigma} = -g\eta - \frac{1}{2}(B^2 + |V|^2) \end{cases} \quad (1.35)$$

then  $\partial_t \Phi$  will be the variational solution of the problem

$$\Delta_{x,y}(\partial_t \Phi) = 0, \quad \partial_t \Phi|_{\Sigma} = -g\eta - \frac{1}{2}(B^2 + |V|^2).$$

By uniqueness, we deduce from (1.13) that  $\partial_t \Phi = -Q$ , which completes the proof of Proposition 1.2. Therefore we are left with the proof of (1.35).

Recall that  $\tilde{\psi}(t, x, z) = \chi(z)e^{z(D_x)}\psi(t, x)$ . Moreover by Lemma 1.3 we have